A tömeg eredete a látható világegyetemben

Katz Sándor

ELTE Elméleti Fizikai Tanszék MTA-ELTE "Lendület" rácstérelmélet kutatócsoport

(Borsányi Szabolcs, Stephan Dürr, Fodor Zoltán, Christian Hoelbling, Stefan Krieg, Laurent Lellouch, Thomas Lippert, Antonin Portelli, Szabó Kálmán, Tóth Bálint)

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Katz Sándor lattice QCD

Outline



Introduction

- 2 Contributions to the proton/neutron masses
- Isospin symmetric case
- Isospin splittings



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Motivation

Composition of the Universe

68% dark energy
27% dark matter
5% ordinary matter
Dark energy and dark matter are yet unknown

Ordinary matter (visible Universe):

Standard Model particles and their bound states mostly (in mass) atoms and nuclei 99.9% of the mass comes from protons & neutrons proton, neutron:

not elementary particles, built from quarks their masses are non-trivial, a consequence of the strong, electromagnetic and weak interactions

Big bang nucleosynthesys and nuclei chart

 $\Delta m_N = m_n - m_p = 0.0014 m_p$

 Δm_N too small \rightarrow inverse β decay leaving predominantly neutrons $\Delta m_N \approx 0.05\%$ would already lead to much more *He* and much less *H* \rightarrow stars would not have ignited as they did

 $\Delta m_N > 0.14\% \rightarrow$ much faster beta decay, less neutrons after BBN burning of *H* in stars and synthesis of heavy elements difficult

The whole nuclei chart is based on precise value of Δm_N



The 3 interactions in the Standard Model

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Strong (QCD)
   guarks' internal degree of freedom: color
  SU(3) symmetry in color space
   coupling \approx 1
Electromagnetic
   among charged particles
   U(1) symmetry after the electroweak breaking
   coupling: \approx 1/137
Weak interaction
   original SU(2) \times U(1) symmetry spontaneously broken by the Higgs
   no exact symmetry at low energies
   approximate isospin symmetry is a remnant
  coupling: \approx 10^{-5}
+1: gravity
  coupling: \approx 10^{-40}
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lattice QCD

Origin of mass

Masses of elementary particles

Higgs mechanism:

Higgs-fermion Yukawa couplings (λ_i) lead to masses

 $m_i = \lambda_i v / \sqrt{2}$

This is the contribution of the weak interaction

 $m_e = 0.511 \text{ MeV}$

Quark masses are not uniquely defined. E.g. in $\overline{MS}(2\text{GeV})$ scheme: $m_u = 2.2(1) \text{ MeV}$ $m_d = 4.8(2) \text{ MeV}$

Mass of the *H* atom: $m_H = 940 \text{ MeV}$

Where does the rest come from?

Mass of the proton and neutron

Proton and neutron are not elementary:

 $p = uud \quad n = udd$

Besides the quark masses all energy (strong and EM) contributes

Strong interaction

quarks and gluons confined into the hadrons have huge kinetic and interaction energiesThis is the dominant contribution (99%)

Electromagnetic interaction

Charged particles have a static electric field which stores energy

Weak interaction

responsible for the quark masses through the Higgs mechanism negligible beyond that

Determination of the proton/neutron masses

Approximate solution

Isospin symmetry:
All properties of the *u*, *d* quarks are identical their charges are neglected
good approximate symmetry but leads to identical *m_n* and *m_p*Only QCD and the quark masses contribute

More detailed analysis

isospin symmetry is not exact $m_u \neq m_d$ and $Q_u = 2/3$ $Q_d = -1/3$ QCD + QED + quark masses contribute estimate:

quark masses: $\Delta m_N \approx 2.6 \text{ MeV}$

electrostatic energy of the proton: $\Delta m_N \approx -0.9 \text{ MeV}$

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Quantum Chromodynamics (QCD)

QCD: Currently the best known theory to describe the strong interaction. SU(3) gauge theory with fermions in fundamental representation. Fundamental degrees of freedom:

- gluons: A^{a}_{μ} , a = 1, ..., 8
- quarks: ψ , $3(color) \times 4(spin) \times 6(flavor)$ components

pure gauge part

where

fermionic part

 $\mathcal{L}_{\text{QCD}} = -\frac{1}{4} F^{a}_{\mu\nu} F^{a\mu\nu} + \overline{\psi} (i D_{\mu} \gamma^{\mu} - m) \psi,$

Quantum Chromodynamics (2)

 \mathcal{L}_{QCD} is invariant under local gauge transformations:

$$egin{aligned} &\mathcal{A}_{\mu}'(x)=G(x)\mathcal{A}_{\mu}(x)G(x)^{\dagger}-rac{i}{g}\left(\partial_{\mu}G(x)
ight)G(x)^{\dagger}\ &\psi'(x)=G(x)\psi(x)\ &\overline{\psi}'(x)=\overline{\psi}(x)G^{\dagger}(x) \end{aligned}$$

Only gauge invariant quantities are physical.

Properties of QCD:

• Asymptotic freedom:

Coupling constant $g \to 0$ when energy scale $\mu \to \infty$.

 \implies Perturbation theory can be used at high energies.

Confinement:

Coupling constant is large at low energies.

 \implies Nonperturbative methods are required.

Quantum Chromodynamics (3)

Quantization using Feynman path integral:

 $\langle 0 | T[\mathcal{O}_{1}(x_{1})\cdots\mathcal{O}_{n}(x_{n})] | 0 \rangle = \frac{\int [d\psi] [d\overline{\psi}] [dA_{\mu}] \mathcal{O}_{1}(x_{1})\cdots\mathcal{O}_{n}(x_{n}) e^{iS[\psi,\overline{\psi},A_{\mu}]}}{\int [d\psi] [d\overline{\psi}] [dA_{\mu}] e^{iS[\psi,\overline{\psi},A_{\mu}]}}$

 e^{iS} oscillates \longrightarrow hard to evaluate integrals. Wick rotation: $t \rightarrow -it$ analytic continuation to Euclidean spacetime. $\implies e^{iS} \longrightarrow e^{-S_{\mathsf{E}}}$, where

$$\mathcal{S}_{\mathsf{E}} = \int \mathrm{d}^4 x \ \mathcal{L}_{\mathsf{E}} = \int \mathrm{d}^4 x \left[\frac{1}{4} \mathcal{F}^a_{\mu
u} \mathcal{F}^a_{\mu
u} + \overline{\psi} (\mathcal{D}_\mu \gamma^\mu + m) \psi \right]$$

positive definite Euclidean action.

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Lattice regularization



take a finite spacetime volume (*TV*) and discretize with a lattice gluons live on links, quarks live on sites path integral \rightarrow finite ($\mathcal{O}(10^9)$) number of integrals quarks integrated analytically: $S \rightarrow S_{eff}$ numerical Monte Carlo integration possible for the gauge fields

Importance sampling

generate configurations with $\exp(-S_{eff})$ weight

Particle masses

two point functions: $\langle O(0)O(t)\rangle \xrightarrow[t \to \infty]{} const \cdot exp(-mt)$

High performance computing required



Juqueen system in Jülich

- 458752 cores
- 6 Pflop/s peak performance
- 2 Pflop/s sustained performance for Monte-Carlo codes

GPU cluster at the Eötvös University



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GPU cluster at ELTE

- 387072 cores
- 1.1 Pflop/s peak performance
- 78 Tflop/s sustained performance

Isospin symmetric case



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masses extracted from various two-point functions correlation length \rightarrow hadron masses perfect agreement with experiments

The challenge of computing $m_n - m_p$

Unprecedented precision is required $\Delta m_N/m_N = 0.14\% \rightarrow$ sub-permil precision is needed to get a high significance on Δm_N

 $m_u \neq m_d \rightarrow 1+1+1+1$ flavor lattice calculations are needed \rightarrow algorithmic challenge (Previous QCD calculations were typically 2+1 or 2+1+1 flavors)

Inclusion of QED: no mass gap

- ightarrow power-like finite volume corrections expected
- ightarrow long range photon field may cause large autocorrelations

Overview of full QED-QCD simulations

- Fully dynamical lattice calculation of SU(3) × U(1) gauge theory with non-degenerate u,d,s,c quarks
- Interpolation/extrapolation to physical point (continuum limit, physical quark masses)
- Treat power-like finite volume corrections analytically/numerically
- Potential precision: $\mathcal{O}(\frac{1}{N_c m_b^2}, \alpha^2) \sim 10^{-4}$
- Calculate mass splittings of N, Σ, Ξ as well as charmed hadrons, D, Ξ_{cc} and Coleman-Glashow relation
- Reaching the desired accuracy via a blind analysis (hard)
- Determine QCD/QED contributions of splittings

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Zero mode subtraction

The absence of a mass gap may cause divergences at finite volume perturbative momentum sums $\rightarrow 1/k^2$ factors \rightarrow zero mode problematic

Removing a finite number of modes does not change $V \rightarrow \infty$ physics

Advantages of zero mode removal:

 \rightarrow analytic calculation of finite V corrections is possible

 \rightarrow algorithmic speedup

Many possibilities, we use:

(Hayakawa, Uno) all spatial zero modes: $\sum_{\vec{x}} A_{\mu, x_0, \vec{x}} = 0 \quad \forall \mu, x_0$

QED in a finite volume

Calculate 1 loop self energy of charged particles in finite and infinite volumes

For a point-like particle

$$\Delta\Sigma(p,L) = \left[\sum_{k} - \int \frac{d^{4}k}{(2\pi)^{4}}\right] \sigma(k,p)$$

zero mode removal: $\sum_{k} \equiv \frac{1}{TL^{3}} \sum_{k_{0}} \sum_{\vec{k}\neq 0}$

Finite V correction to the pole mass can be calculated Result for a spin half particle:

$$m(L) \sim_{L \to +\infty} m \left\{ 1 - q^2 \alpha \left[\frac{\kappa}{2mL} \left(1 + \frac{2}{mL} \right) - \frac{3\pi}{(mL)^3} \right] \right\}$$
with $\kappa = 2.837297(1)$

$$\kappa_{L} \sim 10^{-10} \text{ for } \kappa_{L} \sim 10$$

Composite particles

In QCD charged hadrons are not point-like

previous results have to be extended

QED Ward identities \rightarrow first two orders universal:

$$m(T,L) \underset{T,L \to +\infty}{\sim} m\left\{1 - q^2 \alpha \frac{\kappa}{2mL} \left[1 + \frac{2}{mL}\right] + \mathcal{O}(\frac{\alpha}{L^3})\right\}$$

for scalars and spin half fermions

Form factors (e.g. charge radius) enter at $\mathcal{O}(\frac{\alpha}{13})$ level.

Strategy:

include analytic corrections for the two universal orders fit coefficient of $1/L^3$

 $1/L^3$ in many cases negligible, only significant for mesons

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Finite V dependence of the kaon mass



Neutral kaon shows no volume dependence Volume dependence of the K splitting is perfectly described $1/L^3$ order is significant

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Determining the isospin splittings

two sources of isospin violation: electromagnetism & $m_u \neq m_d$

we work at various elecromagnetic couplings and quark masses renormalized coupling defined at hadronic scales

 \Rightarrow only linear term in α

 $\delta m = m_d - m_u$ is very small \Rightarrow linear term is also enough for δm

 $\Rightarrow \Delta M_X = F_X(M_{\pi^+}, M_{K^0}, M_{D^0}, L, a) \cdot \alpha + G_X(M_{\pi^+}, M_{K^0}, M_{D^0}, a) \cdot \Delta M_K^2$

 ΔM_K^2 : QED-like *L* dependence: absorbed in $F_X(..., L, a)$ charged particle masses: corrected for universal finite-size effect non-universal effects starting with $1/L^3$ are allowed in the QED part

alternative procedure: use ΔM_{Σ} (less precise) instead of ΔM_{K}^{2}

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Separating QED and QCD effects

if α or $(m_u - m_d)$ vanishes QED or QCD parts disappear separation: $\Delta M_X = \Delta_{\text{QED}} M_X + \Delta_{\text{QCD}} M_X$

it is sufficient to decompose the kaon mass squared difference use $\Delta M_X = F_X(M_{\pi^+}, M_{K^0}, M_{D^0}, L, a) \cdot \alpha + G_X(M_{\pi^+}, M_{K^0}, M_{D^0}, a) \cdot \Delta M_K^2$ separation ambigous: depends on the choice of scheme for $m_u - m_d$

Suggested separation:

use Σ^+ and Σ^- baryons they have the same charge squared and same spin leading order: mass difference comes from quark masses only

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Coleman-Glashow relation (1961)

- suggestion: $\Delta_{CG} = \Delta M_N \Delta M_{\Sigma} + \Delta M_{\Xi}$ is close to zero
- remember the quark compositions, in which it cancels indeed: M(ddu) - M(uud) - M(sdd) + M(suu) + M(ssd) - M(ssu) = 0
- determine the leading order terms in the α and δm expansion
- for α =0 a complete quark exchange symmetry $\Delta_{CG} \propto (m_s - m_d)(m_s - m_u)(m_d - m_u)$
- for $\alpha > 0$ remains a $d \leftrightarrow s$ symmetry, thus $\Delta_{CG} \propto \alpha (m_s m_d)$

the Coleman-Glashow relation is satisfied to high accuracy $\Delta_{CG} = 0.00(11)(06) \; \text{MeV}$

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Isospin splittings



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	splitting [MeV]	QCD [MeV]	QED [MeV]
ΔN=n-p	1.51(16)(23)	2.52(17)(24)	-1.00(07)(14)
$\Delta \Sigma = \Sigma^{-} - \Sigma^{+}$	8.09(16)(11)	8.09(16)(11)	0
$\Delta \Xi = \Xi^{-} - \Xi^{0}$	6.66(11)(09)	5.53(17)(17)	1.14(16)(09)
$\Delta D = D^{\pm} - D^{0}$	4.68(10)(13)	2.54(08)(10)	2.14(11)(07)
$\Delta \Xi_{cc} = \Xi_{cc}^{++} - \Xi_{cc}^{+}$	2.16(11)(17)	-2.53(11)(06)	4.69(10)(17)
$\Delta_{CG} = \Delta N - \Delta \Sigma + \Delta \Xi$	0.00(11)(06)	-0.00(13)(05)	0.00(06)(02)

Dependence on fundamental parameters

How strongly does Δm_N depend on α and the quark masses?



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Summary

- Mass of the visible universe: 99.9% protons and neutrons
- Their masses almost identical, but the tiny difference has cosmological consequences
- Higgs is responsible for 1% of the mass
- The rest is given mostly by QCD and a small contribution by QED
- Both masses have been determined using lattice calculations
- The mass difference is a result of two competing effects (quark masses and charges)
- Good agreement with experiments + QCD/QED contributions given separately
- 4 further splittings determined, 2 of them are predictions