

Skálainvariáns dinamika táguló térben

Somfai Ellák

University of Warwick, Wigner FK SZFI

Adnan Ali

Stefan Grosskinsky

Robin Ball

University of Warwick

Introduction

Scale invariant trajectories – random walk, (fractional) Brownian motion

- $dX \sim dt^\gamma$ (eg. $\gamma = 1/2$)
eventually slower than ballistic
- in 1D: all trajectories meet (with probability 1)

Our question:

- what if space expands faster than $\langle |dX| \rangle$?

Introduction – Expanding space

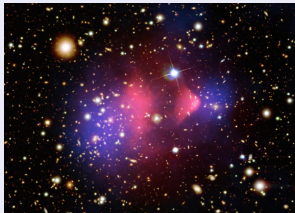
Cosmology



physicsforme.wordpress.com

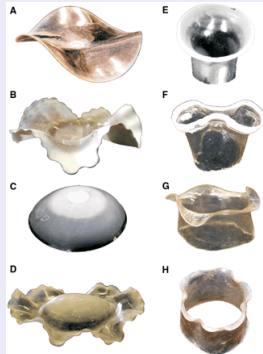
Introduction – Expanding space

Cosmology



physicsforme.wordpress.com

Thin sheets



Klein, Efrati, Sharon, Science (2007)

Introduction – Expanding space

Cosmology



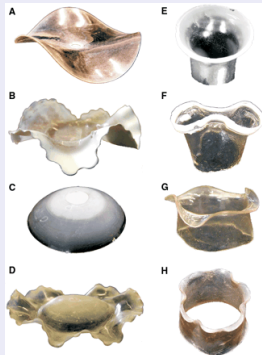
physicsforme.wordpress.com

Growing substrate



stanford.edu/group/brainsinsilicon

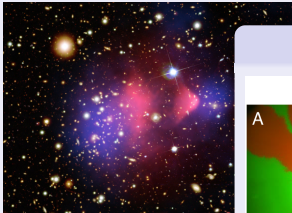
Thin sheets



Klein, Efrati, Sharon, Science (2007)

Introduction – Expanding space

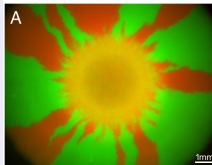
Cosmology



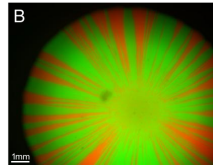
physicsforme.wordp

Indirect: domain walls

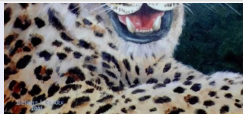
E. coli



S. cerevisiae

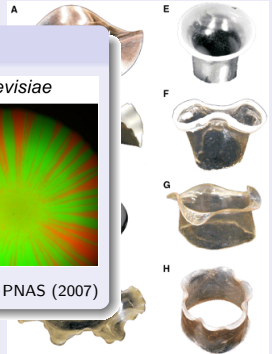


Hallatschek, Hersen, Ramanathan, Nelson; PNAS (2007)



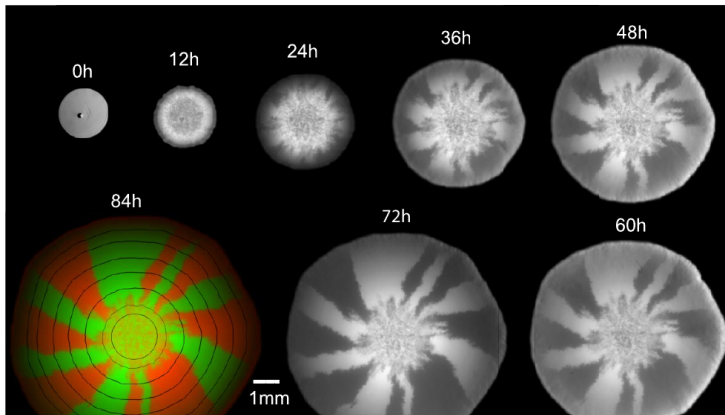
stanford.edu/group/brainsinsilicon

Thin sheets



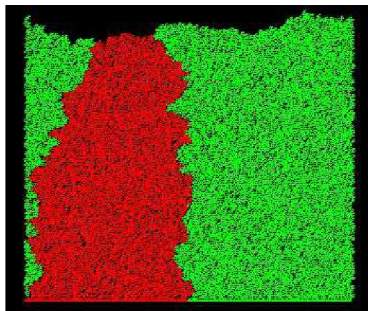
Klein, Efrati, Sharon, Science (2007)

Introduction – Genetic drift and range expansion



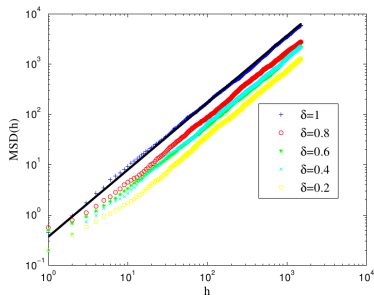
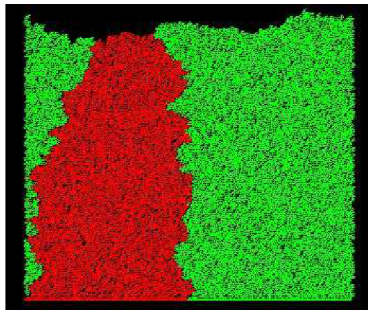
Hallatschek, Hersen, Ramanathan, Nelson; PNAS (2007)

Domain boundaries



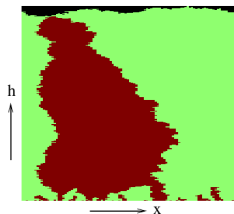
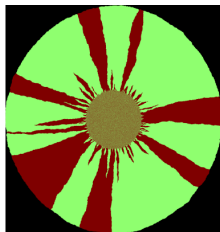
- domain boundaries grow perpendicular to surface
 X_h is superdiffusive due to surface roughness

Domain boundaries



- domain boundaries grow perpendicular to surface
 X_h is superdiffusive due to surface roughness
- $M(h) := \langle X_h^2 \rangle \approx \sigma^2 h^{2\gamma}$, $\gamma = 2/3$

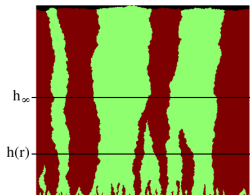
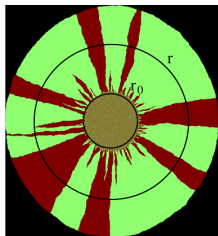
Range expansion



- fixed, finite size geometry: fixation (absorbing state in finite time)
- range expansion: promotes diversity and segregation

Mapping

- understand radial growth with help of same dynamics in fixed size rectangular space

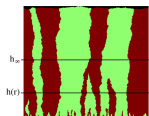


- polar-like coordinate transformation:

linear: $0 \leq X_h < L, \quad 0 \leq h < \infty$

radial: $0 \leq Y_r < 2\pi r, \quad r_0 \leq r < \infty, \quad 2\pi r_0 = L$

Mapping



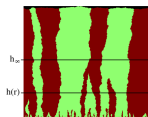
- polar-like coordinate transformation:

linear: $0 \leq X_h < L, \quad 0 \leq h < \infty$

radial: $0 \leq Y_r < 2\pi r, \quad r_0 \leq r < \infty, \quad 2\pi r_0 = L$

$$Y_r = \frac{r}{r_0} X_h$$

Mapping



- polar-like coordinate transformation:

linear: $0 \leq X_h < L, \quad 0 \leq h < \infty$

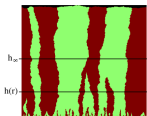
radial: $0 \leq Y_r < 2\pi r, \quad r_0 \leq r < \infty, \quad 2\pi r_0 = L$

$$Y_r = \frac{r}{r_0} X_h$$

- increment:

$$dY_r = \frac{dr}{r_0} X_h + \frac{r}{r_0} dX_h = Y_r \frac{dr}{r} + d\tilde{Y}_r, \quad \frac{r}{r_0} dX_h = d\tilde{Y}_r$$

Mapping



- polar-like coordinate transformation:

$$\text{linear: } 0 \leq X_h < L, \quad 0 \leq h < \infty$$

$$\text{radial: } 0 \leq Y_r < 2\pi r, \quad r_0 \leq r < \infty, \quad 2\pi r_0 = L$$

$$Y_r = \frac{r}{r_0} X_h$$

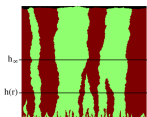
- increment:

$$dY_r = \frac{dr}{r_0} X_h + \frac{r}{r_0} dX_h = Y_r \frac{dr}{r} + d\tilde{Y}_r, \quad \frac{r}{r_0} dX_h = d\tilde{Y}_r$$

- preserving local structure:

$$dX_h \sim (dh)^\gamma, \quad d\tilde{Y}_r \sim (dr)^\gamma \quad \begin{array}{ll} \gamma = 1/2 & \text{diffusive fluctuations} \\ \gamma = 2/3 & \text{KPZ domain boundaries} \end{array}$$

Mapping



- polar-like coordinate transformation:

$$\text{linear: } 0 \leq X_h < L, \quad 0 \leq h < \infty$$

$$\text{radial: } 0 \leq Y_r < 2\pi r, \quad r_0 \leq r < \infty, \quad 2\pi r_0 = L$$

$$Y_r = \frac{r}{r_0} X_h$$

- increment:

$$dY_r = \frac{dr}{r_0} X_h + \frac{r}{r_0} dX_h = Y_r \frac{dr}{r} + d\tilde{Y}_r, \quad \frac{r}{r_0} dX_h = d\tilde{Y}_r$$

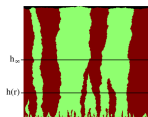
- preserving local structure:

$$dX_h \sim (dh)^\gamma, \quad d\tilde{Y}_r \sim (dr)^\gamma \quad \begin{array}{ll} \gamma = 1/2 & \text{diffusive fluctuations} \\ \gamma = 2/3 & \text{KPZ domain boundaries} \end{array}$$

-

$$\frac{dh}{dr} = \left(\frac{dX_h}{d\tilde{Y}_r} \right)^{1/\gamma} = \left(\frac{r_0}{r} \right)^{1/\gamma}$$

Mapping



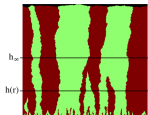
- integrating differential equation:

$$h(r) = \begin{cases} r_0 \frac{\gamma}{1-\gamma} \left(1 - (r_0/r)^{\frac{1-\gamma}{\gamma}} \right) & , \gamma \neq 1 \\ r_0 \ln(r/r_0) & , \gamma = 1 \end{cases}$$

- local interaction:

$$\left\{ \frac{r_0}{r} Y_r \right\} \stackrel{\text{dist.}}{=} \{ X_{h(r)} \} \quad \text{for all } r \geq r_0$$

Mapping



- integrating differential equation:

$$h(r) = \begin{cases} r_0 \frac{\gamma}{1-\gamma} \left(1 - (r_0/r)^{\frac{1-\gamma}{\gamma}}\right) & , \gamma \neq 1 \\ r_0 \ln(r/r_0) & , \gamma = 1 \end{cases}$$

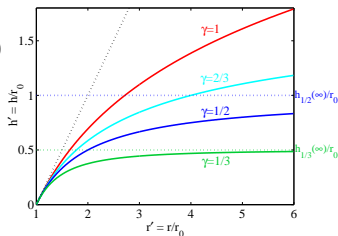
- local interaction:

$$\left\{ \frac{r_0}{r} Y_r \right\} \stackrel{\text{dist.}}{=} \{X_{h(r)}\} \quad \text{for all } r \geq r_0$$

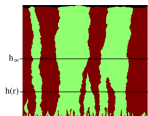
- properties:

$$h(r) \approx r - r_0 \quad \text{for } r \text{ close to } r_0$$

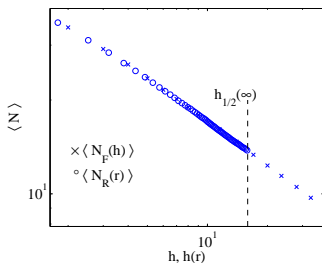
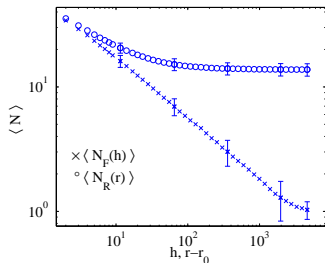
$$h_\gamma(\infty) = \lim_{r \rightarrow \infty} h(r) = \frac{\gamma}{1-\gamma} r_0 < \infty$$



Results



- number of domains:



Results

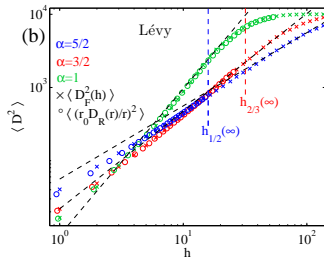
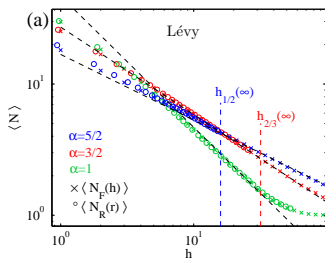
Different processes:

- Lévy flight:
jump size distribution: $\mathbb{P}(X_{h+} - X_h = x) \sim C|x|^{-(1+\alpha)}$ ($\alpha > 0$)
Markovian
 $\gamma = \max\{1/\alpha, 1/2\}$
- fractional Brownian motion:
covariances: $\langle X_{h+\Delta h} X_h \rangle \sim (h + \Delta h)^{2\gamma} + h^{2\gamma} - (\Delta h)^{2\gamma}$
non-Markovian

Results

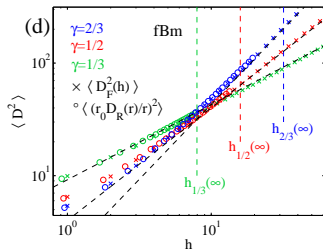
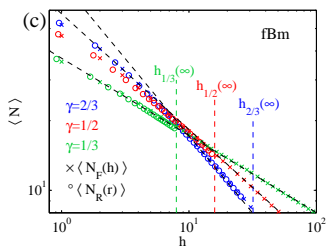
Lévy flight ($\gamma = \max\{1/\alpha, 1/2\}$) + absorption:
number of surviving trajectories N ,

mean square distance: $D_F(h)^2 = \sum_{i=1}^{N(h)} \left(X_h^{(i+1)} - X_h^{(i)} \right)^2$



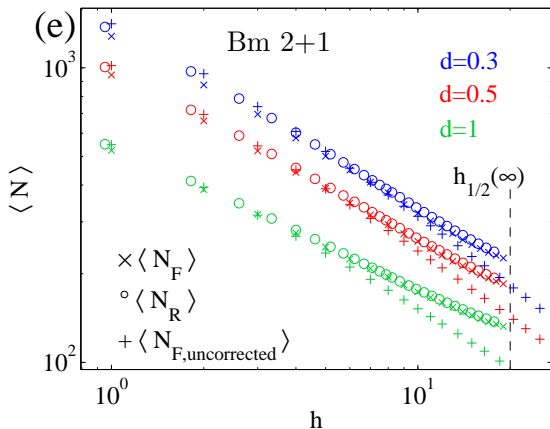
Results

fractional Brownian motion:



Results – finite particle size

fixed size d in **expanding** space corresponds to decreasing size $\frac{r_0}{r(h)}d$ in **fixed space**

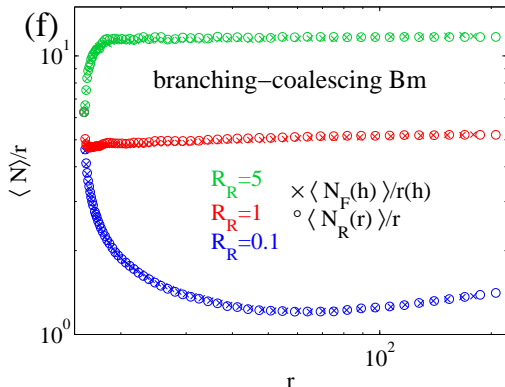


Results – branching-coalescing Brownian motion

branching rate R_R in expanding space vs.

branching rate R_F in fixed space:

$$\frac{R_R}{R_F} = \frac{\Delta_R(dr)/dr}{\Delta_F(dh)/dh} = \frac{dh}{dr} = \left(\frac{r_0}{r}\right)^{1/\gamma}$$

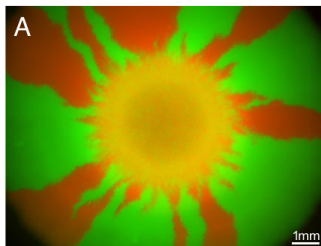


Conclusions – mapping

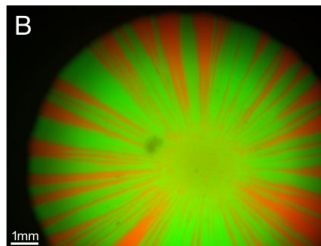
- locally scale invariant growth processes mapped from homogenously expanding space to fixed space
- can be used eg. to handle 2D radial growth, **asymptotic state** in radial growth corresponds to **finite time** in fixed space
- can be extended to include:
 - ▶ finite particle size,
 - ▶ branching-coalescing,
 - ▶ higher dimensions etc.

Bacteria vs. yeast

E. coli

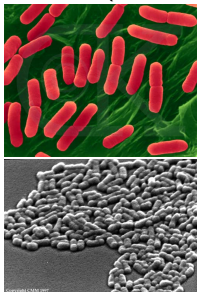


S. cerevisiae



Hallatschek, Hersen, Ramanathan, Nelson; PNAS (2007)

Bacteria (*E. Coli*)



emc.maricopa.edu

Yeast (*S. cerevisiae*)

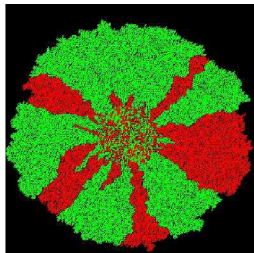
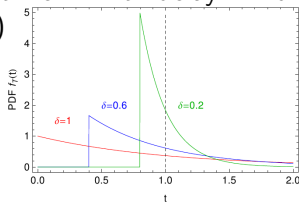


visualphotos.com

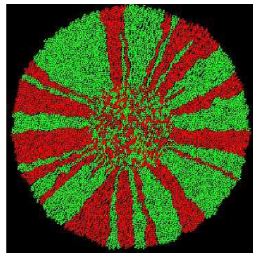
- ignore many biological details (shape, growth direction, etc.)
- consider: correlations due to reproduction time
- consider: overall geometry

Model

- off-lattice Eden growth model
- one-parameter family, $\delta \in [0, 1]$: reproduction time T with delay $1 - \delta$
distribution: $T \sim 1 - \delta + \text{Exp}(1/\delta)$
normalized average: $\langle T \rangle = 1$
variation coefficient: $\sigma(T)/\langle T \rangle = \delta$



$\delta = 1$



$\delta = 0.2$

Surface – KPZ scaling

- on large scale:

$$\partial_t h(x, t) = v_0 + \nu \partial_x^2 h + \frac{\lambda}{2} (\partial_x h)^2 + D \eta(x, t)$$

- scaling:

$$x \rightarrow x' = bx, \quad t \rightarrow t' = b^z t, \quad h \rightarrow h' = b^\alpha h$$

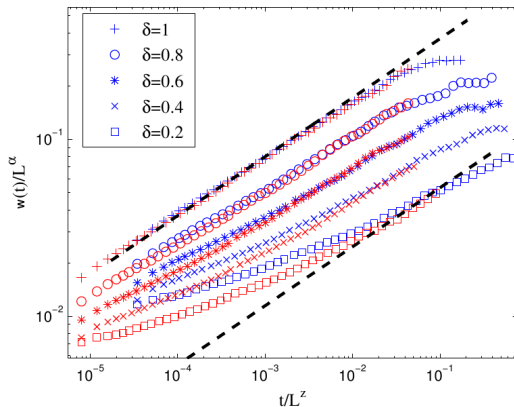
$$\text{statistical scale invariance: } h(x, t) \sim h'(x', t')$$

$$\text{Family-Vicsek scaling: } w(L, t) := \sqrt{\langle (h - \langle h \rangle_x)^2 \rangle_x} \sim L^\alpha f(t/L^z)$$

$$\text{ahol } f(u) \sim \begin{cases} u^\beta, & \text{if } u \ll 1 \\ \text{const}, & \text{if } u \gg 1 \end{cases}$$

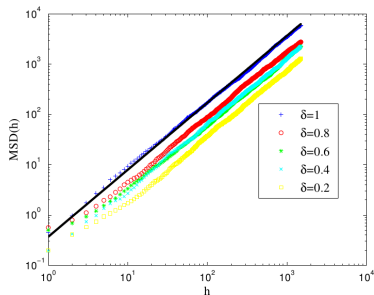
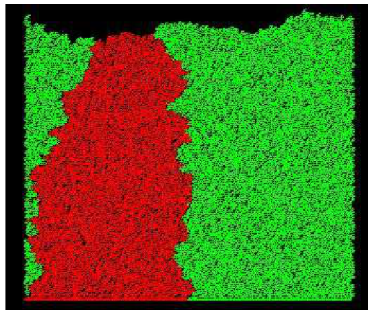
- exponents: $\alpha + z = 2, \quad z = \alpha/\beta$
1D: $\alpha = 1/2, \quad \beta = 1/3, \quad z = 3/2.$

Surface – KPZ scaling



$$w(L, t) \sim \begin{cases} t^{1/3}, & \text{for } t \ll L^{3/2} \\ L^{1/2}, & \text{for } t \gg L^{3/2} \end{cases}$$

Domain boundaries



- domain boundaries grow perpendicular to surface
- X_h is superdiffusive due to surface roughness
- $M(h) := \langle X_h^2 \rangle \approx \sigma^2 h^{2\gamma}, \quad \gamma = 2/3$

Correlations

Partial synchronization leads to intrinsic vertical correlations.

- $N(t)$ growth events with height Δh_i : $h_{N(t)} = \sum_{i=1}^{N(t)} \Delta h_i$
- $\text{var}[h_{N(t)}] = \langle \Delta h_i \rangle^2 \text{var}[N(t)] + \langle N(t) \rangle \text{var}[\Delta h_i]$
 $= t \langle \Delta h_i \rangle^2 (\delta^2 + \epsilon^2) \stackrel{!}{=} O(1)$

where correlation coefficient due to geometric effects:

$$\epsilon = \sqrt{\text{var}[\Delta h_i]} / \langle \Delta h_i \rangle$$

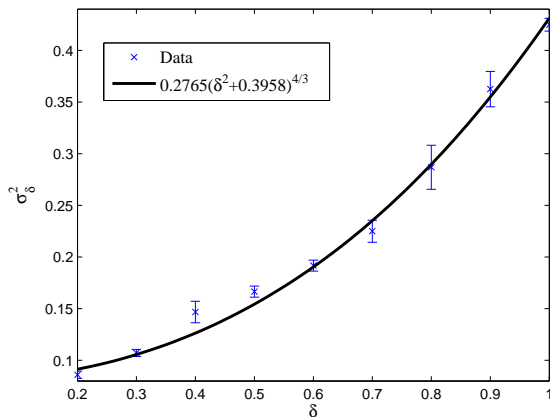
- intrinsic vertical correlation: $\tau \sim \frac{1}{\delta^2 + \epsilon^2}$
lateral correlation length: $\xi_{\parallel}(t) \sim (t/\tau)^{1/2}$
- mean square displacement

$$M(h) := \langle [X(h) - X(0)]^2 \rangle \approx \sigma_{\delta}^2 h^{2\gamma} \sim \xi_{\parallel}^2(h)$$

$$\gamma = 2/3 \quad \text{and} \quad \sigma_{\delta}^2 \propto (\delta^2 + \epsilon^2)^{4/3}$$

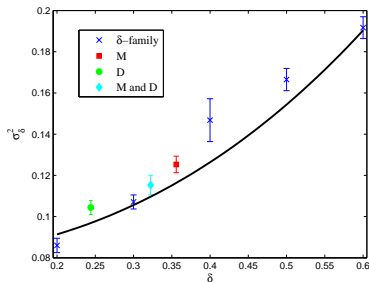
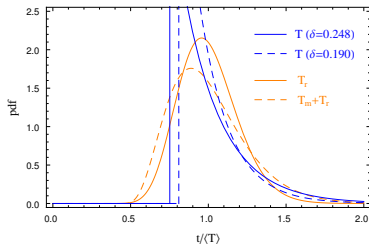
Correlations

Prefactor of mean square displacement: $\sigma_\delta^2 \propto (\delta^2 + \epsilon^2)^{4/3}$



S. cerevisiae (yeast)

Replace δ -family with realistic reproduction times.



Conclusions – partially synchronized growth

- extension of off-lattice Eden model
reproduction time has variation coefficient δ
- stays in KPZ universality class
- changes in patterns are due to changing prefactors – quantified
- works for realistic reproduction times (*S. cerevisiae*)